Comparative Analysis of Plasma Probe Diagnostics Techniques

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Acknowledgements:

This work was supported by the DOE OFES (Contract No DE-SC0001939) The authors are thankful to their colleagues for sharing with us their data

68th Gaseous Electronics Conference,

9th International Conference on Reactive Plasmas, 33rd Symposium on Plasma Processing October 12 – 16, 2015, Hawaii Convention Center, Honolulu, USA

Introduction

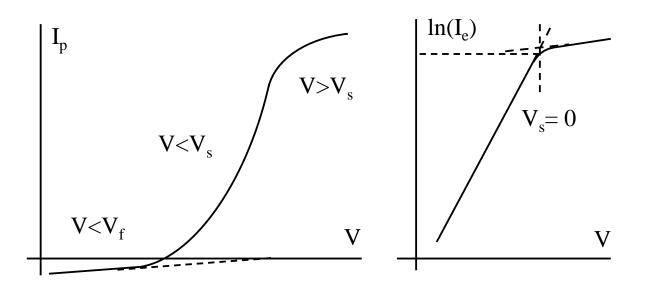
"There is no plasma diagnostics method other than probe diagnostics where the danger of incorrect measurements and erroneous interpretation of results are so great."

L. Schott, in *Plasma Diagnostics*, editor W. Lochte-Holtgreven, Amsterdam,1968

It was true then, it is even more true today, when plasmas are more complicated and the measurement techniques are more sophisticated

Comparison of plasma parameters obtained with different probe techniques (Langmuir, ion current and EEDF) is given in this talk for collisionless probes at B = 0

Electron part of the probe I/V (Classical Langmuir technique)



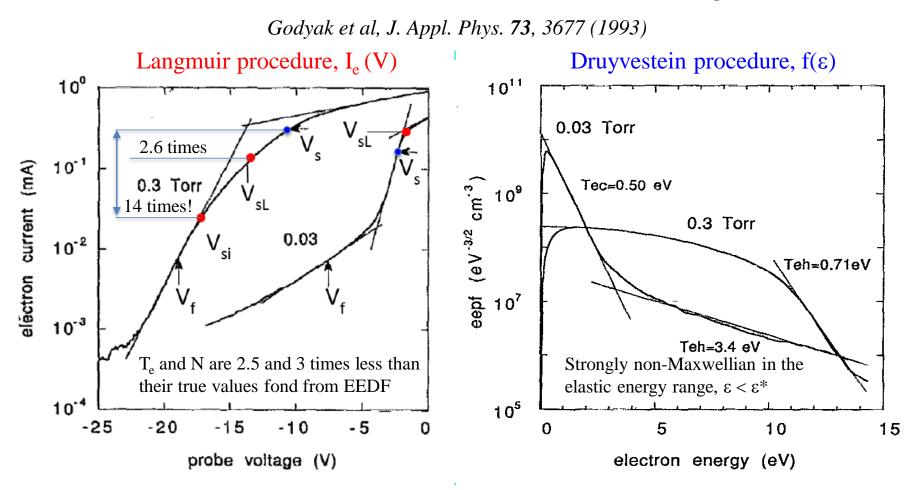
Large probe gives a nice looking probe I/V, but prone to plasma perturbation

Main problems in using Langmuir procedure:

- 1. Assumption of Maxwellian EEDF \longrightarrow errors in T_e and N
- 2. Uncertainty in plasma potential V \longrightarrow error in N

3. Arbitrariness in the ion current approximation \longrightarrow error in T_e for high energy electrons ($\varepsilon > |eV_f|$)

Probe characteristics and EEDFs in Argon CCP



Deviation from Maxwellian EEDF for high energy electrons ($\varepsilon > \varepsilon^*$, ε_i , ε_w) is typical for gas discharge plasmas. Strong non-equilibrium for bulk electrons ($\varepsilon < \varepsilon^*$, ε_i , ε_w) in DC and RF discharges in Ramsauer gases makes the <u>classical Langmuir procedure</u> inadequate for diagnostics of such plasmas

Plasma parameters for Ar CCP at 30 and 300 mTorr

30 mTorr (bi-	Maxwellian) _I	300 mTorr (Druyvesteyn-like)				
Langmuir	EEDF	Langmuir	EEDF			
$T_{e} = 0.73 \text{ eV}$	$T_{e} = 0.67 eV$	$T_{e} = 1.37$	$T_{e} = 3.4 \text{ eV}$			
N=5.9x10 ⁹ cm ⁻³	$N=4.4 \times 10^9 \text{ cm}^{-3}$	N=4.5x10 ⁹ cm ⁻³	N=2.9x10 ⁹ cm ⁻³			

The temperatures of cold and fast electrons in bi-Maxwellian plasma found from the Langmuir procedure and those found from the EEDF are different!

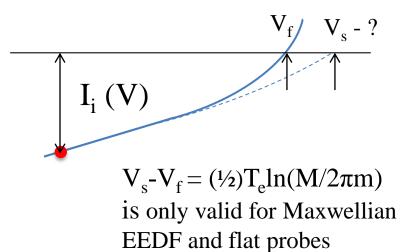
$$T_{ec} = 0.73 \text{ eV}$$
 and $T_{he} = 4.2 \text{ eV}$ are found as $T_e^{-1} = d \ln I_e / dV$
 $T_{ec} = 0.50 \text{ eV}$ and $T_{he} = 3.4 \text{ eV}$ are found as $T_e^{-1} = d \ln f_e / dV$

Inferring the plasma parameters from the ion part of the probe characteristic

 $T_e = [d(I - I_i)/I_i dV]^{-1} @ V = V_f$

N is found from one of theories: Radial, or Orbital motion, Kagan-Perrel or Laframboise. Which one to use ?

Numerous studies showed that plasma parameters found from the ion part of the probe characteristic can be in error, up to an order of magnitude comparing to those found from the electron part of the probe characteristic! $\downarrow \downarrow \downarrow \downarrow \downarrow \downarrow$

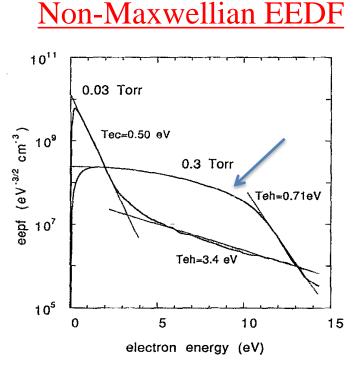


<u>Warning!</u>

The inferred T_e is formed by the fast electrons of the EEDF, T_{eh} . The last could be larger, or less than T_{eeff} defined by the bulk electrons.

Possible reasons for "ion density" failure? <u>Unrealistic assumptions for I_i(V) theories:</u>

- Maxwellian EEDF,
- No ion collisions $(\lambda_i << \lambda_e)$,
- One-dimensional probe models



- No ambipolar flow $(v_s \ge v_{amb} > v_{Ti})$
- Uncertainty in V_s
- Unknown ion temperature

As shown by Vasil'eva (*High Temp. 12, 409, 1974*), I_i and $\lambda_D \sim [T_{es}]^{1/2}$, while V_f $\sim T_{eh}$, where T_{es} is the electron screen temperature waited by T_{ec}.

@ 0.3Torr, $T_{es} = 15 \text{ eV}$, $T_{eh} = 0.71 \text{ eV}$, and $T_e = 3.4 \text{eV}$.

 T_e at $V = V_f$ is an order of magnitude lower than T_{es} that defines the Bohm velocity, u_B and the sheath collecting area around the probe. Combination of these effects may result in order of magnitude error in N_i. **Ion collisions** Generally, $\lambda_i \ll \lambda_e$, and collisional regime for ion onsets at much lower gas pressure than that for electrons. For the orbital ion theory:

 $R_0 = a_p (eV/T_i)^{1/2} \ll \lambda_i$, @ $a_p = 0.05$ mm, $T_i = 0.03$ eV, V = 60 V, the collisionless regime in Ar gas occurs at p << 15 mTorr. I. e. OML theory is not applicable for the range of gas pressure where the theory usually used (p > 2-3 mTorr).

1-D not satisfied

R₀ and S are comparable to the probe length; L = 5 mm

 $I_i \sim V^{1/2}$ for cylinder. and $I_i \sim V$ for sphere - ???

 $R_0 = 2.24$ mm and L/2 = 2.5 mm; they are equal, while applicability of the cylindrical model requires $R_0 \ll L/2$

<u>Ambipolar drift</u> In gas discharge plasmas, except of plasma center, $v_{Ti} < v_i < v_B = (T_e/M)^{1/2}$, while all ion current theories (used in probe Diagnostics) assume $v_i = 0$, or $v_i = v_{Ti}$. Ion directed motion can be accounted for, but the ion velocity number and its direction are *a priory* unknown.

Uncertainty and erroneous estimation of the plasma potential

EEDF measurement with probe I/V differentiation

Druyvestein (1931) showed that $d^2I_e/dV^2 \propto EEPF$, f(ϵ), while EEDF, F(ϵ) $\propto \epsilon^{1/2}f(\epsilon)$ gives plasma parameters and process rates

- Applicable to arbitrary isotropic F(ε)
- Applicability limitations are the same as for Langmuir probe
- Accurate V_s measurement and no ion current effect
- Plasma parameters, rate of e-collisional processes and transport coefficients can be readily found as integrals of the measured F(ε)
- Since low temperature plasmas as a rule are non-Maxwellian, the EEDF measurement is recommended as a reliable probe diagnostics for laboratory and processing plasmas

See recent review on EEDF measurement Godyak and Demidov, J. Phys. D; Appl. Phys. **44**,233001, 2011

Ratio of the measured plasma density to that found by integration of the measured EEDF

First	Year of	Gas/pressure	Electron	Ion	Ion	Hairpin	Interfero-	Cut-off
Author	publication		Part of	orbital	radial	probe	meter	probe
			I/V	theory	theory			
Godyak	1993	Ar 30 mT	1.34	2.5				
Godyak	1993	Ar 300 mT	0.38/0.07	3.3				
Sudit	1994	He 40 mT	0.85	9 !!!	0.25			
Piejak	2004	Ar 3-50 mT				1.2-1.5	1.3-1.6	
Ki	2005	Ar 7-22 mT		2.6-3.2		1.5		1.05
Iza	2006	Ar 1mT		3	0.3			
Iza	2006	Ar 10 mT		4	0.45			
Iza	2006	Ar 100 mT		2 !	0.14 !!!			

Godyak et al, J. Appl. Phys. **73** 3657, 1993 Sudit and Woods, J. Appl. Phys. **76**, 4488, 1994 Piejak et al, J. Appl. Phys. **95**, 3785 (2004) Ki and Chung, Korean Phys. Soc. **50**, 329, 2005 Iza and Lee J. Vac. Sci. Technol. A **24**, 1366, 2006

Large data discrepancy found from ion and electron probe currents.

More reliable are data obtained with minimal number of assumption (EEDF and ω_{0e})

What makes a reliable EEDF measurement?

- EEDF measurements yield meaningful results only when they contain accurate information about electrons in both elastic and inelastic energy range; ($\epsilon < \epsilon^* \rightarrow T_e \& N$) and ($\epsilon > \epsilon^* \rightarrow \epsilon$ excitation & ionization). This requires high energy resolution ($\Delta \epsilon < 0.5 T_e$) and large dynamic range of the measurements, (50-70) db undistorted by noise.
- Measurements in processing chambers have to be unaffected by the wide spectrum of RF and low frequency plasma potential and the probe and chamber surface contamination.
- Real time EEDF monitoring allows to notice impediments in measurements and mitigate them.

Majority of published EEDF data measured by self-made and commercial instruments are obviously distorted missing essential information. That indicates either lack of the probe diagnostics skills or the instrument deficiency, or frequently both.

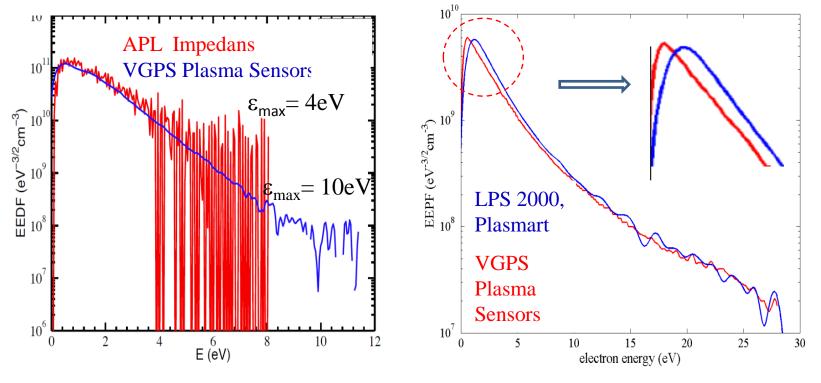
EEDF measured with different instruments in similar plasmas Plasma Sensors VGPS Impedance ALP 10¹⁵ Ar ICP, 0.3; 1 **10**¹³ and 10 mTorr Ar ICP, 2.2 mTorr 6.78 MHz, 50 W 10¹² Lost information EEPF (eV³²².3 about low and high 10¹³ Low energy peak in 1.011 energy electrons EEDF is typical in a eepf (eV ^{-3/2} cm ⁻³ low pressure ICP 10¹⁰ dominated by anomalous skin effect 10⁹ 10¹¹ 10⁸ 300 mT 100 10⁷ 10 10⁹ 0.3 $\epsilon^* \epsilon_i$ 10⁶ 30 20 40 10 0 10 20 30 40 50 0 Electron energy (eV) electron energy (eV)

Gahan et al, PSST 17, 035026 (2008)

Godyak et al, PSST 11, 525 (2002)

EEDF measurements in the same noisy plasma with different commercial instruments

(Dynamic range and energy and resolution)



PEGASES. Ar 10 mTorr, 100W. Ecole Polytechnique, France ICP with ferrite core. Ar 1.2 mTorr, 40 W Tsinghua University, China

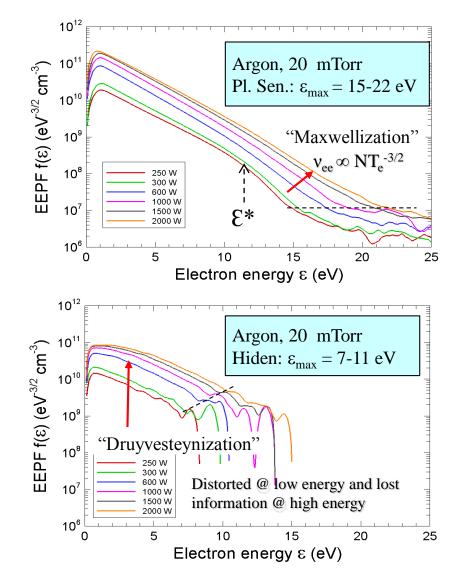
EEDF measurements in a commercial ICP reactor

Comparison of EEPF measured with different commercial probe stations, Espion of Hiden and VGPS of Plasma Sensors

At maximal discharge power of 2 kW, N $\approx 1 \cdot 10^{12}$ cm⁻³, thus the EEPF @ $\epsilon < \epsilon^*$ has to be a Maxwellian one

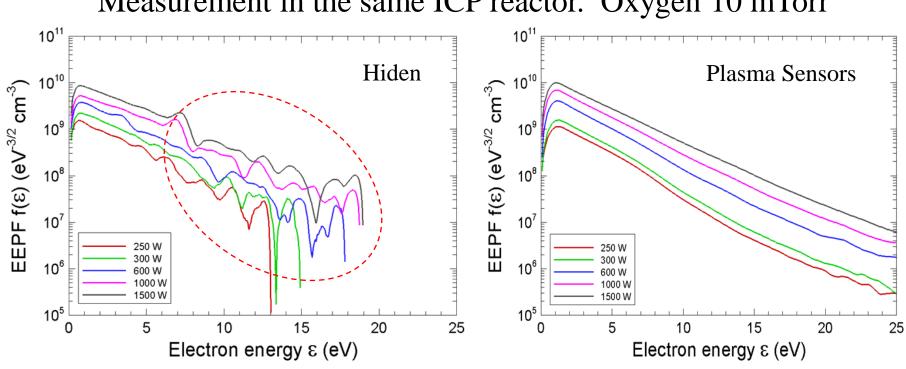
"Druyvesteynization" effect is found in many publications of EEDF measurements made with self-made and commercial probe systems

V. Godyak et al, GEC 2009, Saratoga Springs, NY, USA



Measuring of EEDF in reactors with processing gases

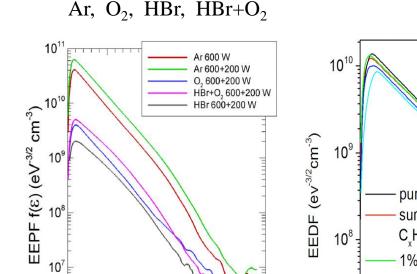
High amplitude and wide spectrum of RF plasma potential harmonics, low-frequency noise, high rate of the probe contamination and poor contact of the plasma to the grounded chamber are major impediments for probe diagnostics in commercial chambers.



Measurement in the same ICP reactor. Oxygen 10 mTorr

V. Godyak et al, GEC 2009, Saratoga Springs, NY, USA

EEDFs measured in plasma reactors with processing gases (High quality measurements with a good instrument)



10⁸ 1% CH, 2% CH, 4% CH,

Electron energy ε (eV)

10

15

20

25

10

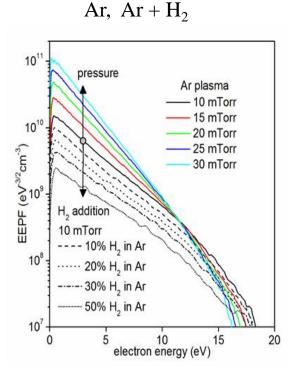
10⁶

0

EEPF in commercial two-inductor ICP in different processing mixtures at 15 mTorr. Mattson Technology, by V. Nagorny

 $Ar + SiH_4$ Ar/CH, plasma 300 W, 10 mTorr pure Ar surface derived C_xH_y flow 10^{7} 15 5 10 20 electron energy (eV) EEPFs in ICP filled with Ar+CH₄ at the condition of strong polymer film deposition. University of Maryland,

by N. Fox-Lyon



EEPF measured in ICP reactor, H_2/CF_4 at 30 mTorr with a polymer layer deposition. University of Maryland by N. Fox-Lyon

Concluding remarks

- The classic Langmuir probe procedure is subjected to errors due to non-Maxwellian EEDF and uncertainty in the plasma potential, it should not be used for plasma diagnostics with non-Maxwellian EEDF at ε <ε*
- Plasma parameters inferred from $I_i(V)$ are significantly (up to an order of magnitude) different from those found from $I_e(V)$ and EEDF. Unrealistic assumptions made in $I_i(V)$ theories are likely suspects for errors in T_e and N values found from $I_i(V)$.
- EEDF probe diagnostics is not confined by those drawbacks, but it requires more accurate instruments and elaborate procedures.
 Deficiencies in both components are visible in vast number of published EEDF measurements.

Refining of "well established" diagnostic technique may lead to new finding

- Having a choice, the most accurate measurement are based on fundamental physical principles with least assumptions and parameters restrictions.
- Comparison of numerous studies presenting the plasma parameters inferred from different probe techniques demonstrates vast discrepancy.
- The EEDF measured by Langmuir probe and the plasma density found from ω_{pe} in the cut-off probe technique are both based on fundamental principles and are potentially most accurate.
- Cut-off probe is immune to the probe contamination and can be made insusceptible to RF potential. But its only output is the plasma density, still within a certain range.
- Accurate EEDF measurements can be performed in wide parameter range revealing most comprehensive plasma information. It allows unambiguous calculation of the plasma parameters and rates of the transport and reaction processes as corresponding integrals of EEDF.

Three levels of probe diagnostics

Plasma parameters are inferred from:

- 1. Ion part of the probe I/V characteristic I_i (double and triple probes). Is notoriously inaccurate (up to an order of magnitude error) due to many unrealistic assumptions in existing $I_i(V)$ theories
- 2. Electron part of the probe I/V (classic Langmuir method). Assumes a Maxwellian EEDF, uncertainty in plasma potential leads to error in the plasma density evaluation
- **3. I/V Differentiation of the probe characteristic** . Results in accurate measurement of the plasma potential and EEDF. Plasma parameters and process rates may be found as corresponding integrals of the measured EEDF

Since EEDFs in gas discharge plasmas are never Maxwellian (at both, $\varepsilon < \varepsilon^*$ and $\varepsilon > \varepsilon^*$), the measurement of EEDF is the only reliable probe diagnostics corresponding to contemporary kinetic level of gas discharge science